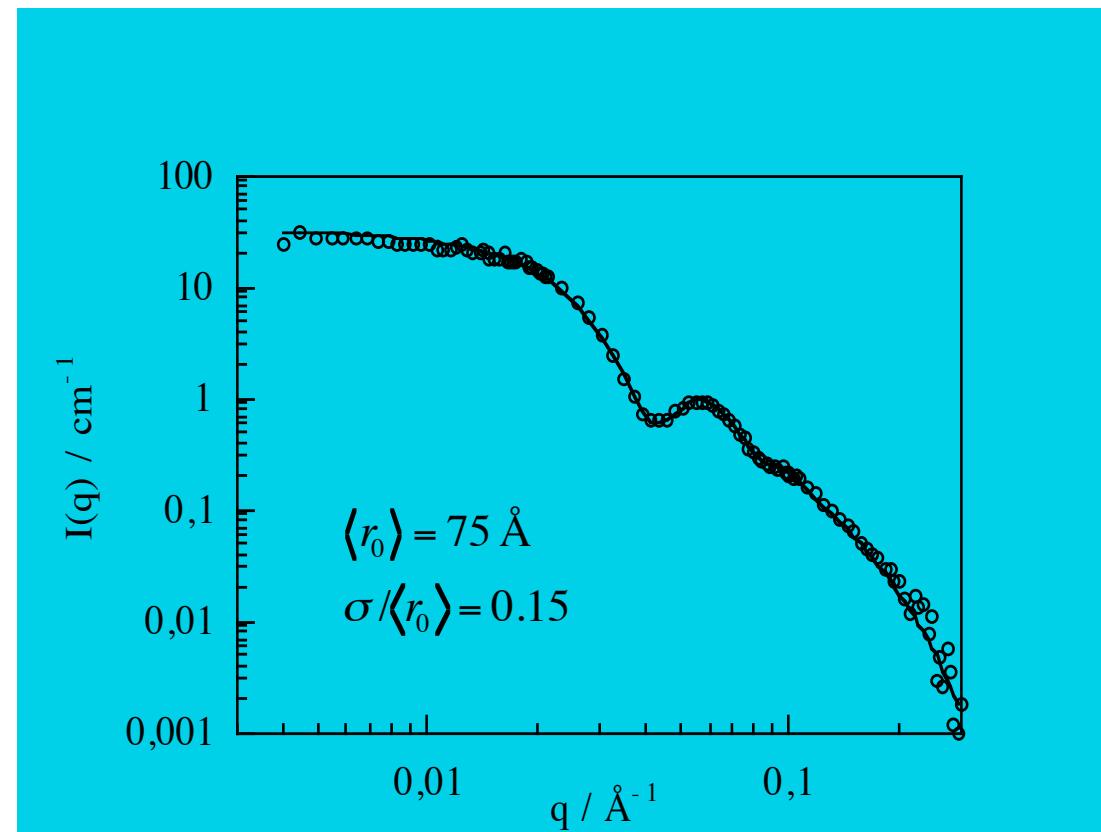
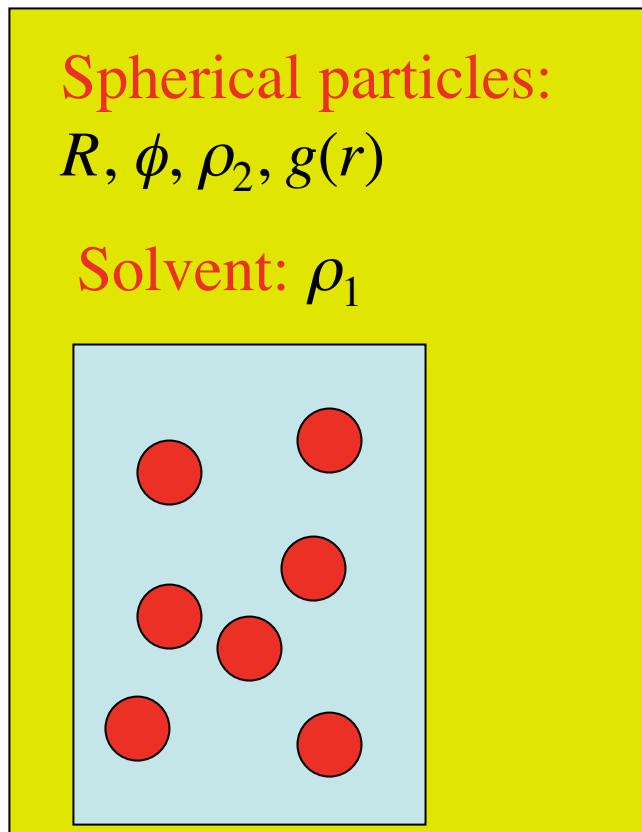
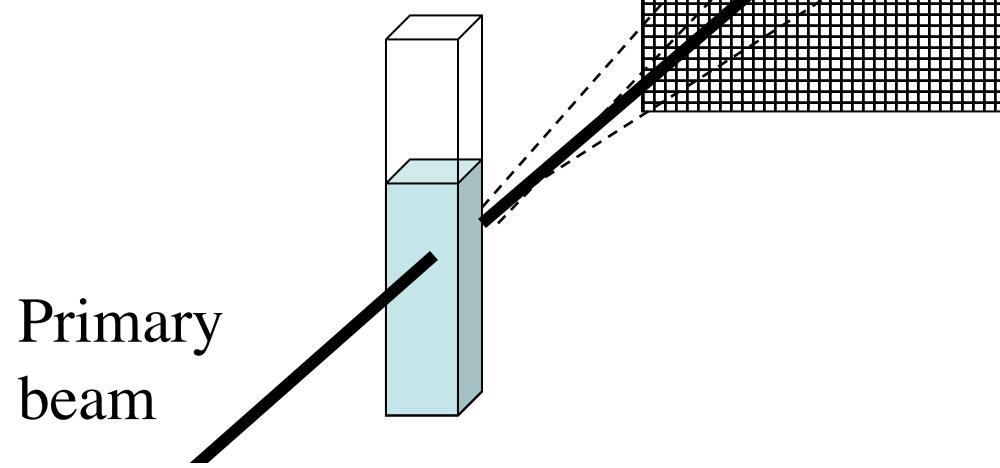


INTRODUCTION TO SMALL ANGLE SCATTERING FROM COLLOIDAL DISPERSIONS

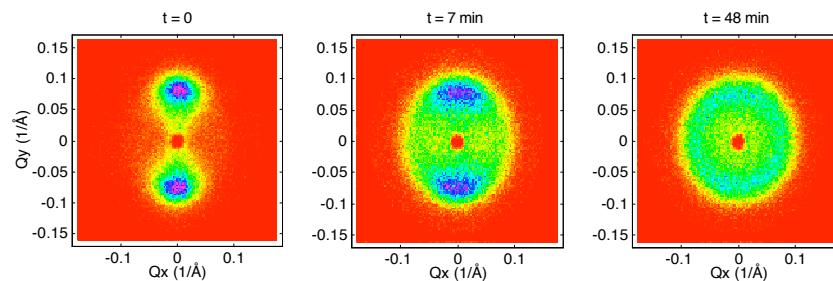
Analysis of small angle scattering data from an isotropic solution (dispersion) of spherical colloidal particles.



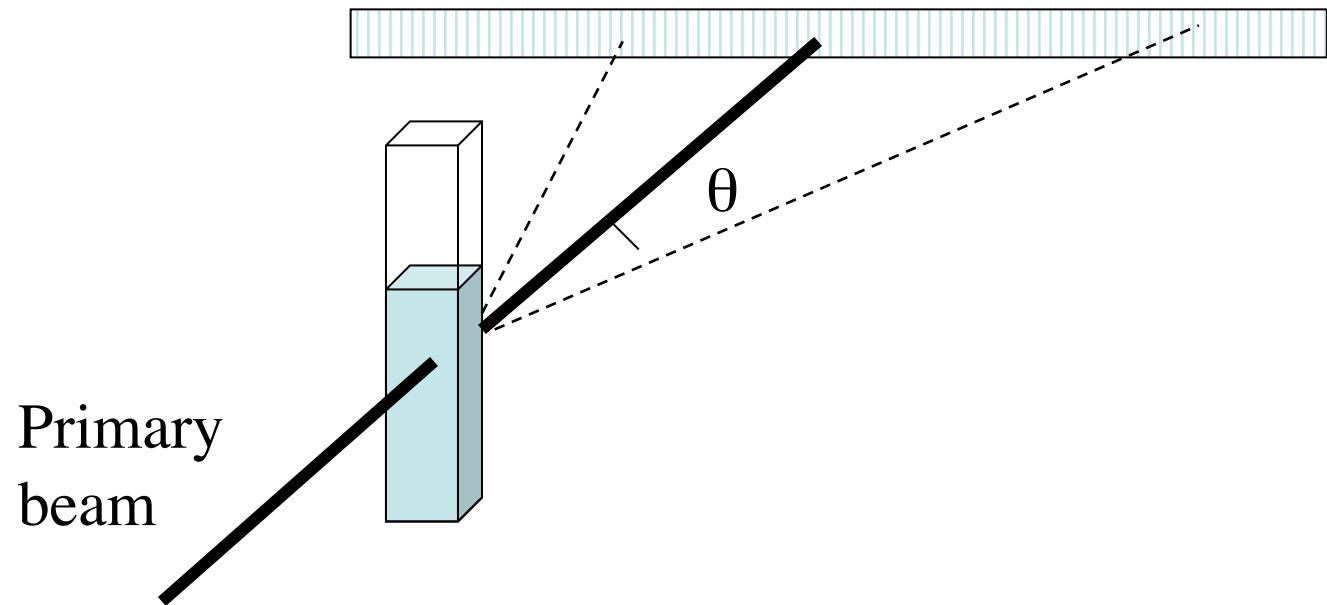
2D detector



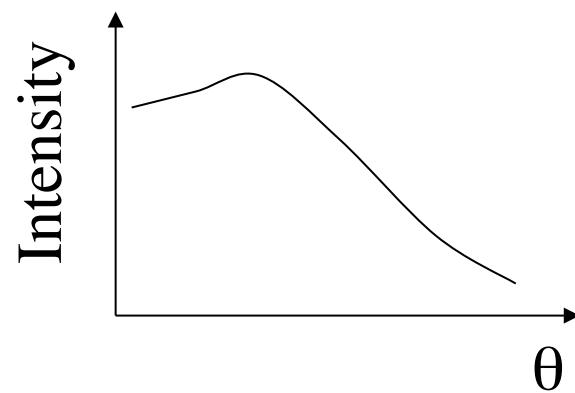
Primary
beam



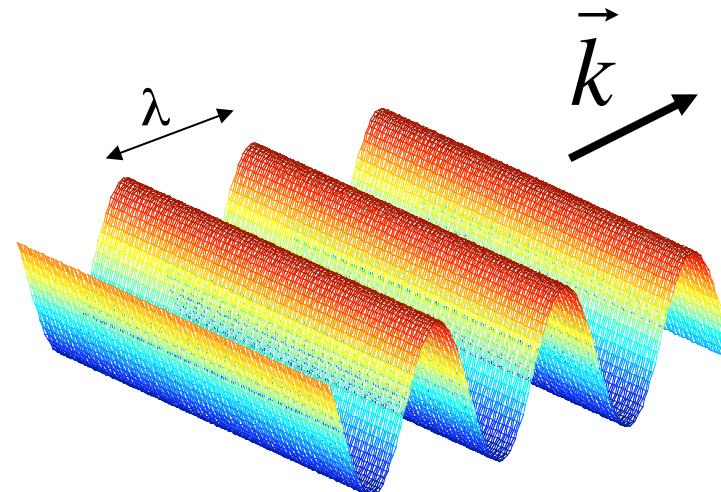
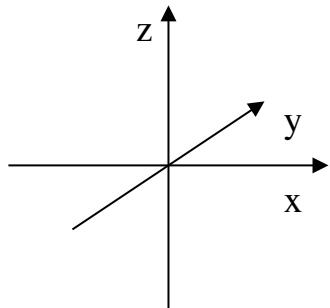
1D detector



Primary
beam



The planar wave



Amplitude:

$$A = A_0 e^{i \vec{k} \cdot \vec{r}}$$

wave vector:

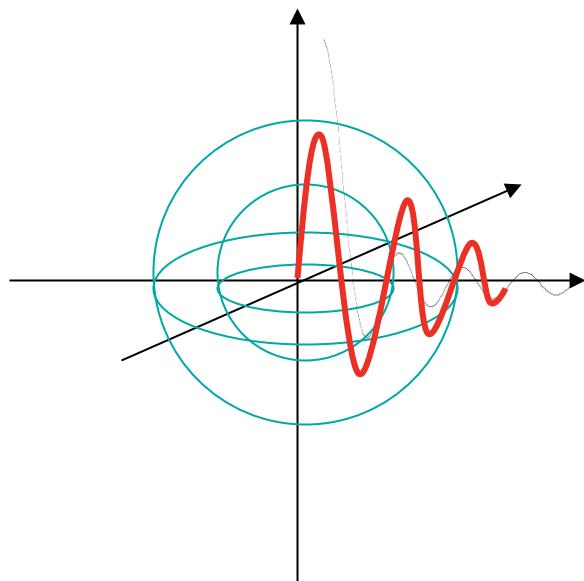
$$|\vec{k}| = \frac{2\pi}{\lambda}$$

Momentum:

$$\vec{p} = \hbar \vec{k}$$

$$|\vec{p}| = \frac{h}{\lambda} \quad (\text{de Broigle})$$

The spherical wave



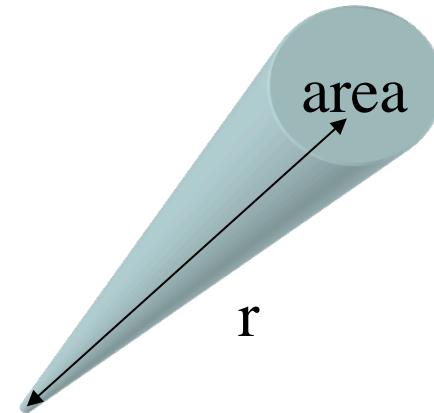
Amplitude:

$$A = A_0 \frac{b}{r} e^{ikr}$$

INTENSITY

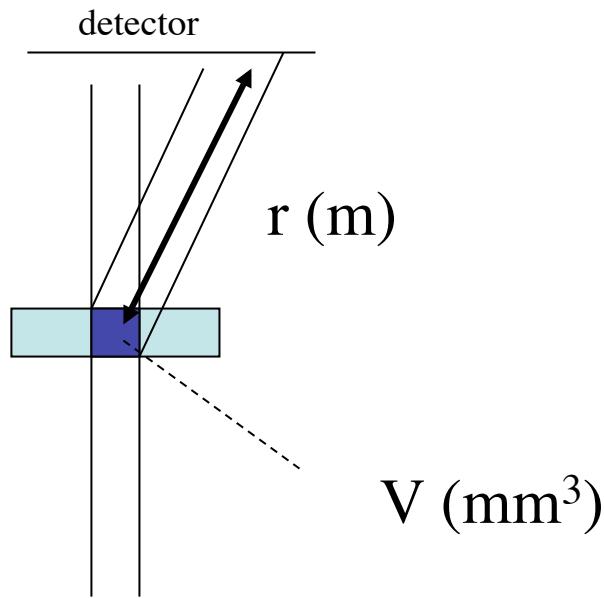
$$\text{Intensity: } I = |A^2|$$

Energy flux / time & area ($\perp \vec{k}$)



”Number of particles (e.g. photons) / time & area”

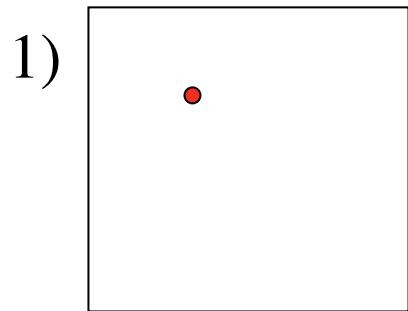
$$Area \sim r^2 \Rightarrow I \sim r^{-2} \Leftrightarrow A \sim r^{-1}$$



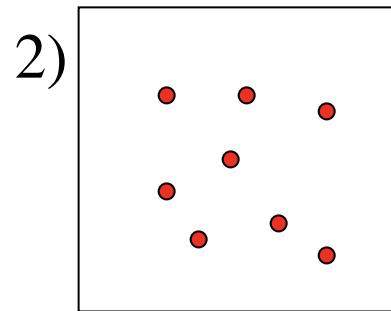
$$I = \frac{I_s}{I_0} \frac{r^2}{V}$$

[length⁻¹] [cm⁻¹]

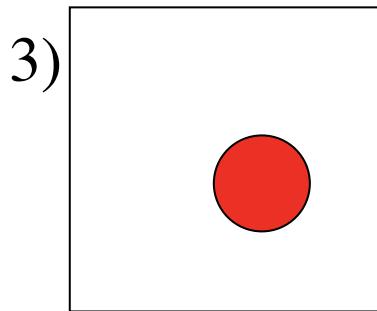
Outline



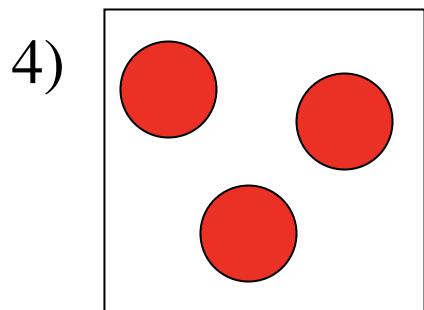
Single scattering centre



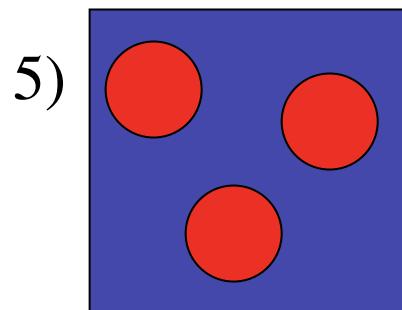
N scattering centra
=> interference



Continuum
Single colloidal particle



N particles
=> interference



+ solvent

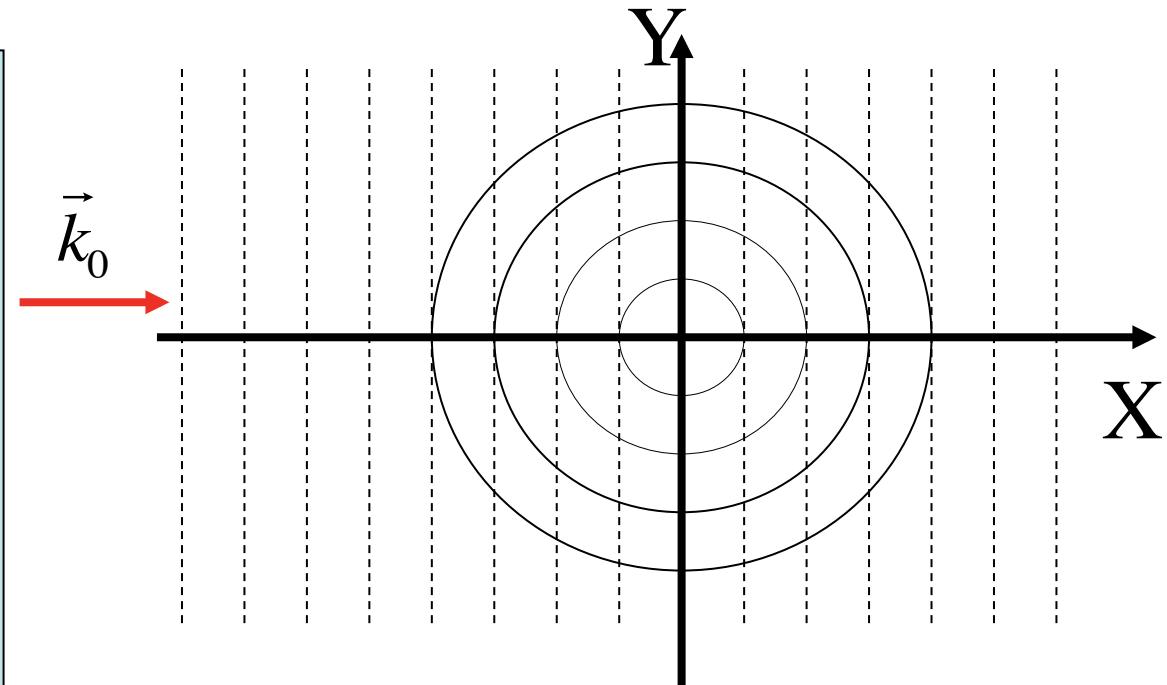
Single scattering centre (single event)

(X-rays: electron, neutrons: atomic nucleus)

$$A_{in} = A_0 e^{i \vec{k}_0 \cdot \vec{r}}$$

$$A_s = A_0 \frac{b}{r} e^{i k_s r}$$

b = "scattering length"



- Coherent scattering: A_{in} and A_s in phase

- Elastic scattering: λ remains unchanged $\Leftrightarrow |\vec{k}_0| = |\vec{k}_s|$

Displaced from the origin - The phase shift

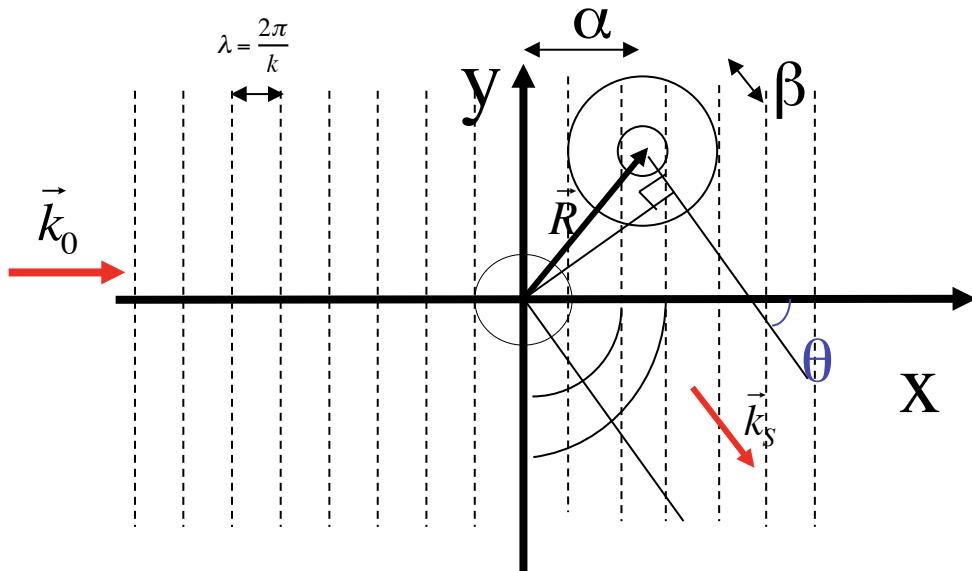
Assume detector distance $r \gg R$

Phase shift =

$$= \frac{2\pi}{\lambda} \alpha + \frac{2\pi}{\lambda} \beta =$$

$$= \vec{k}_0 \cdot \vec{R} + (-\vec{k}_s \cdot \vec{R}) =$$

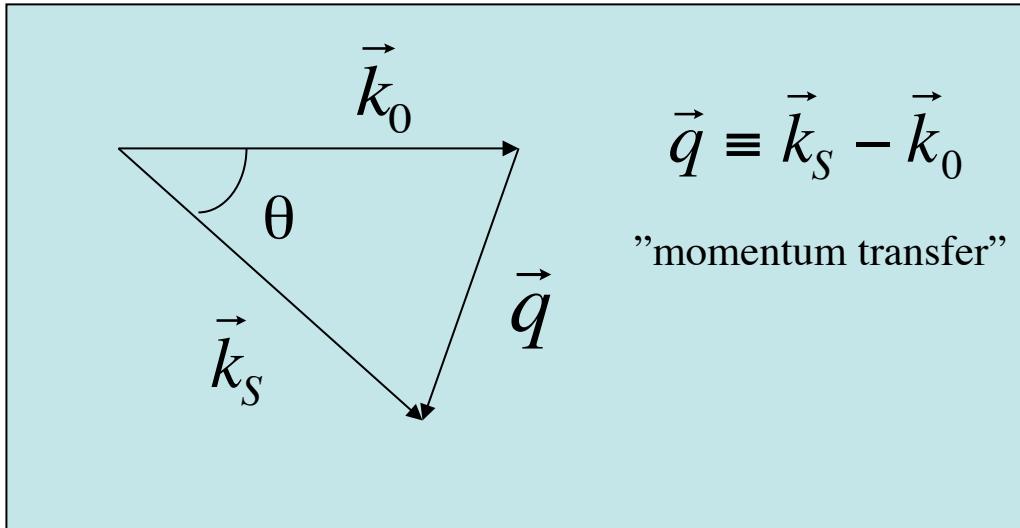
$$= -(\vec{k}_s - \vec{k}_0) \cdot \vec{R} = -\vec{q} \cdot \vec{R}$$



Scattering amplitude:

$$A_S = A_0 \frac{b}{r} e^{i k_s r} e^{-i \vec{q} \cdot \vec{R}}$$

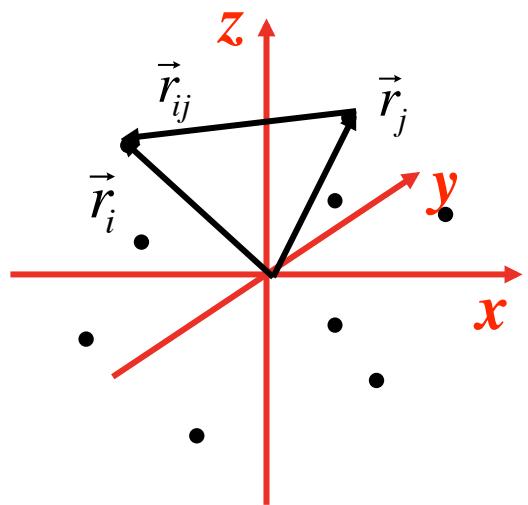
The scattering vector \vec{q}



$$|\vec{k}_S| \equiv k_S = k_0 \quad (\text{elastic})$$
$$\Rightarrow q = \frac{4\pi}{\lambda} \sin\left(\frac{\theta}{2}\right)$$

A diagram illustrating the relationship between the scattering vector \vec{q} and the scattering angle θ . It shows a triangle with vertices representing wave vectors. The top vertex is labeled \vec{q} , the bottom-left vertex is labeled \vec{k}_0 , and the bottom-right vertex is labeled \vec{k}_S . A dashed vertical line connects the top vertex \vec{q} to the hypotenuse of the triangle, representing the projection of \vec{q} onto the plane of the triangle. Below the triangle, the equation $\frac{q}{2} = k_0 \sin\left(\frac{\theta}{2}\right)$ is written.

N scattering centra



Intensity is the square of the amplitude

$$I = \langle |A|^2 \rangle = \frac{\langle |A_S|^2 \rangle}{A_0^2} \frac{r^2}{V}$$

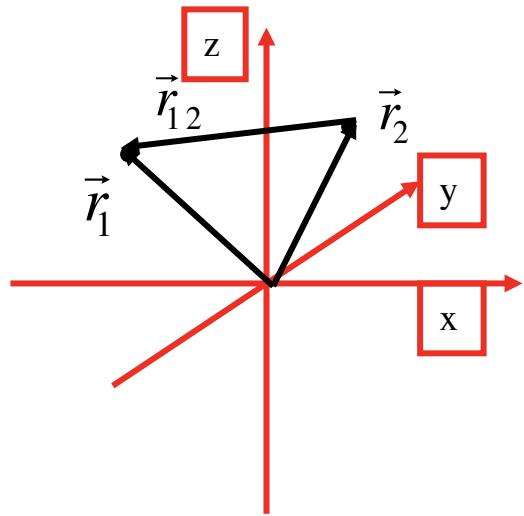
Sum up all the scattered waves:

$$A_S = A_0 \frac{b}{r} e^{ikr} \sum_{i=1}^N e^{-i\vec{q} \cdot \vec{r}_i}$$

$$I = \frac{1}{V} \left\langle b^2 \sum_{i=1}^N \sum_{j=1}^N e^{-i\vec{q} \cdot \vec{r}_{ij}} \right\rangle$$

$$\vec{r}_{ij} \equiv \vec{r}_i - \vec{r}_j$$

The basic idea

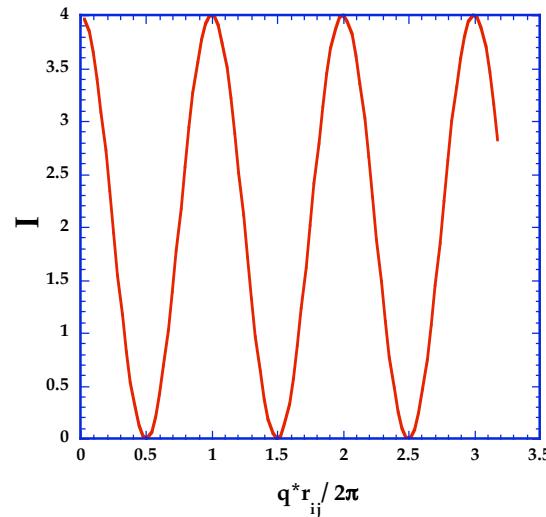


$$N \quad P(q) \quad S(q)$$

$$(=2) \quad (=1)$$

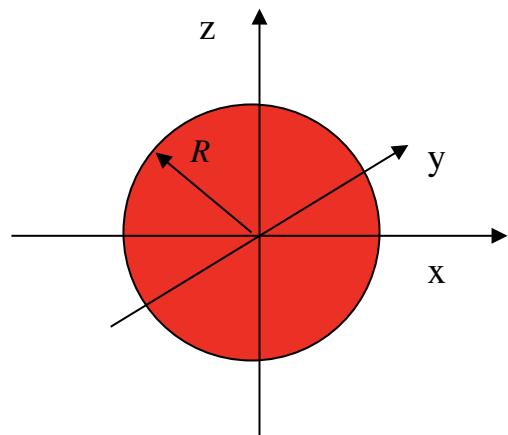
- Maximum for $qr_{12}=n2\pi$.
- $I(q)$ probe correlations at separations $2\pi/q$.

$$\begin{aligned} I(q) &\sim \left(e^{i\vec{q}\cdot\vec{r}_1} + e^{i\vec{q}\cdot\vec{r}_2} \right) \left(e^{-i\vec{q}\cdot\vec{r}_1} + e^{-i\vec{q}\cdot\vec{r}_2} \right) = \\ &= 2 + \left(e^{i\vec{q}\cdot\vec{r}_{12}} + e^{-i\vec{q}\cdot\vec{r}_{12}} \right) = \\ &= 2 \underbrace{\left(1 + \cos(\vec{q} \cdot \vec{r}_{12}) \right)}_{S(q)} \end{aligned}$$



Collect the “atoms” in a sphere

The particle “form factor”

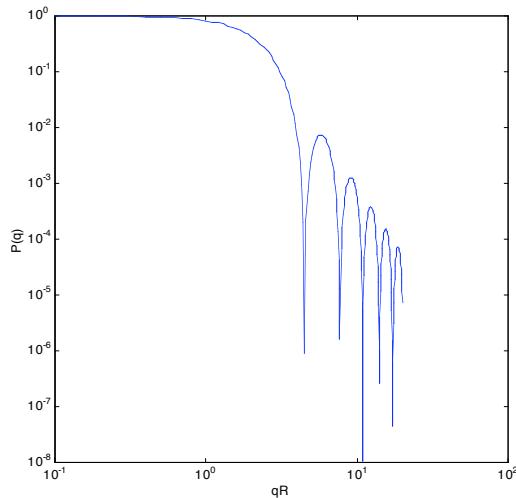


$$\begin{aligned}
 I(\vec{q}) &= |A|^2 = \frac{1}{V} \left| \sum_i b e^{i \vec{q} \cdot \vec{r}_i} \right|^2 \\
 &\approx \frac{1}{V} \left| \int_{V_{sphere}} d\vec{r} \rho(\vec{r}) e^{i \vec{q} \cdot \vec{r}} \right|^2 \\
 &= \frac{1}{V} P(\vec{q})
 \end{aligned}$$

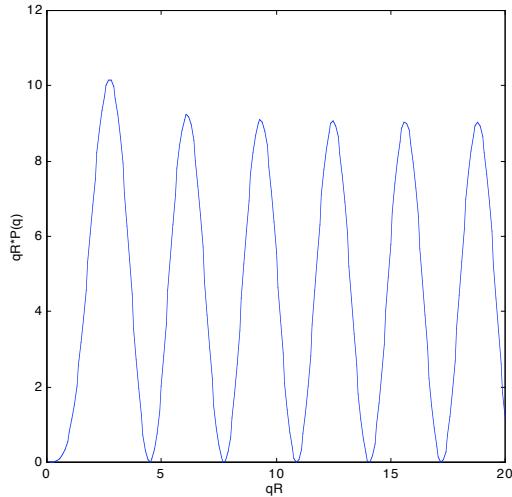
$P(\vec{q})$: “form factor” or single particle scattering function.

$\rho(\vec{r}) = \frac{\partial b}{\partial V}$: “Scattering length density”

The form factor of a sphere.



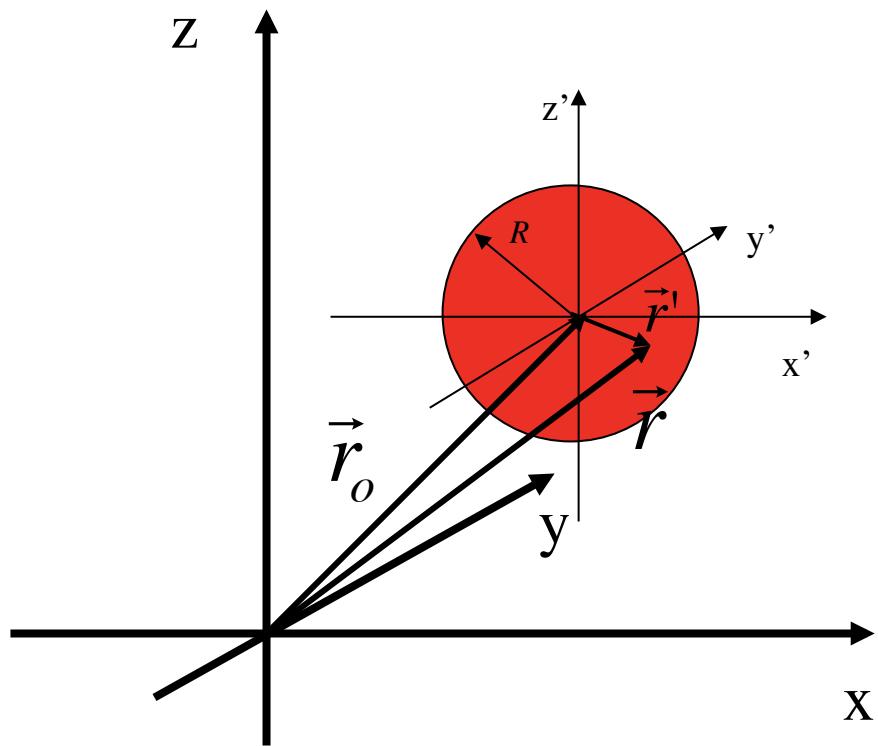
Oscillations damped with q^{-4}



$$\begin{aligned}
 P(q) &= \rho^2 \left| \int_{V_{sphere}} d\vec{r} e^{i\vec{q} \cdot \vec{r}} \right|^2 = \\
 &= \rho^2 \left(3 \frac{4\pi R^3}{3} \frac{\sin\{qR\} - qR \cos\{qR\}}{(qR)^3} \right)^2 = \\
 &= \rho^2 \left(3 \frac{4\pi R^3}{3} \frac{j_1(qR)}{qR} \right)^2
 \end{aligned}$$

- $j_1(x)$: spherical Bessel function
- Zeros at $qR = \tan\{qR\}$

Displacement \vec{r}_o



$[x,y,z]$: “lab-frame”
coordinate system

$[x',y',z']$: local coordinate
system

$$\vec{r} = \vec{r}_o + \vec{r}'$$

$$I(q) = \left| \rho \int_{V_{sphere}} d\vec{r}' e^{i\vec{q} \cdot (\vec{r}' + \vec{r}_o)} \right|^2 = \left| \left(\rho \int d\vec{r}' e^{i\vec{q} \cdot \vec{r}'} \right) e^{i\vec{q} \cdot \vec{r}_o} \right|^2$$

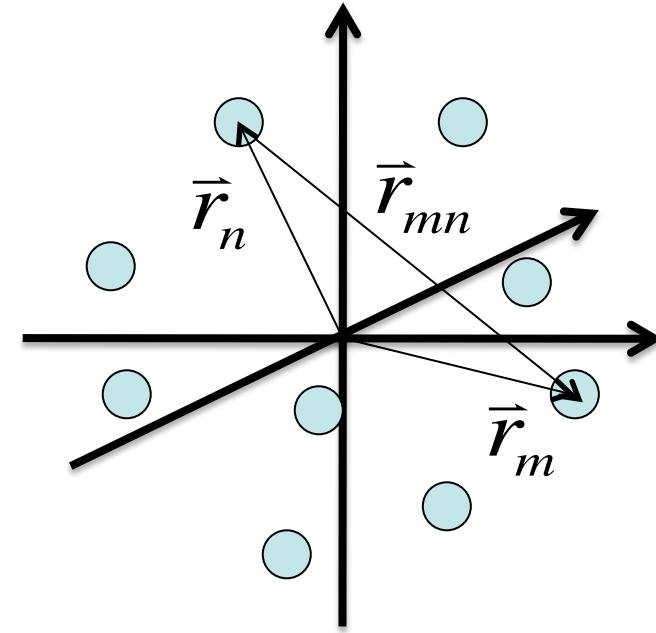
$$I(q) = \frac{1}{V} \left\langle \left| \sum_{m=1}^N \rho \left(\int_{V_s} d\vec{r}' e^{i\vec{q} \cdot \vec{r}'} \right) e^{i\vec{q} \cdot \vec{r}_m} \right|^2 \right\rangle =$$

$$= \frac{1}{V} P(q) \left\langle \sum_{m=1}^N \sum_{n=1}^N e^{i\vec{q} \cdot \vec{r}_{mn}} \right\rangle =$$

$$= \frac{1}{V} P(q) \left\langle N + \sum_{m} \sum_{n \neq m} e^{i\vec{q} \cdot \vec{r}_{mn}} \right\rangle =$$

$$= \frac{N}{V} P(q) \left\langle 1 + \frac{1}{N} \sum_m \sum_{n \neq m} e^{i\vec{q} \cdot \vec{r}_{mn}} \right\rangle =$$

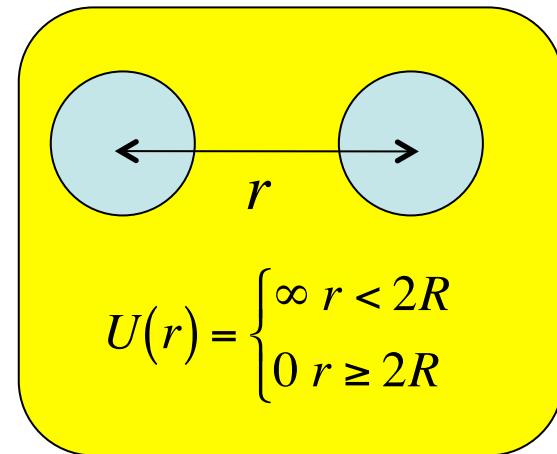
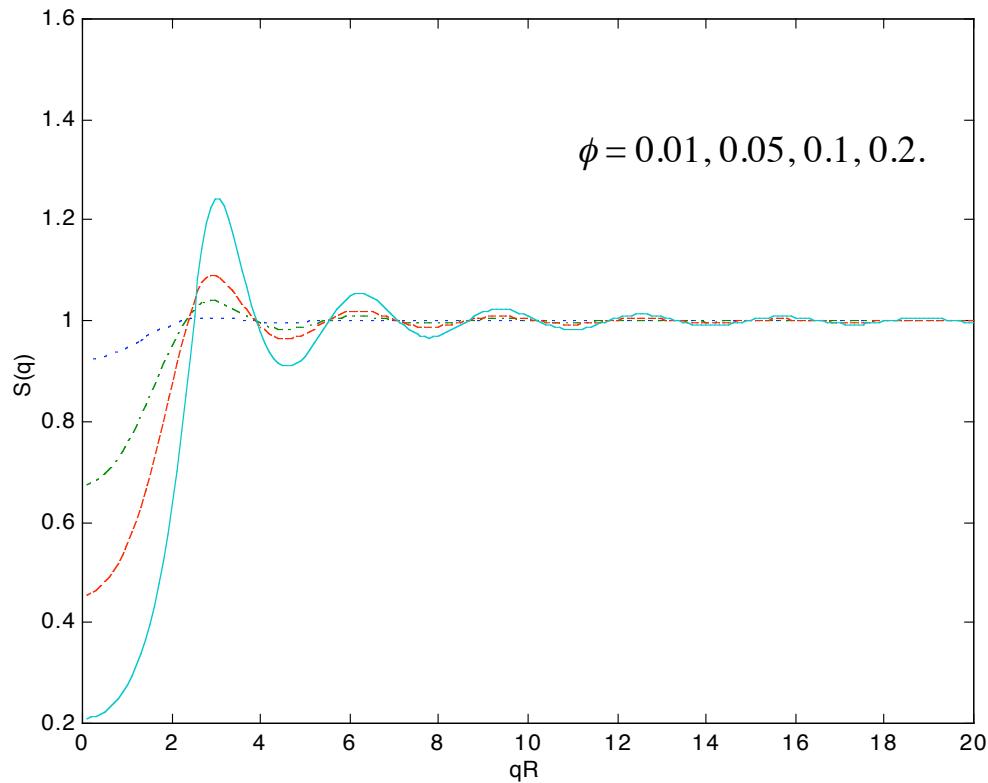
$$= \frac{N}{V} P(q) S(q)$$



$S(q)$: “Structure factor”

Hard sphere structure factor

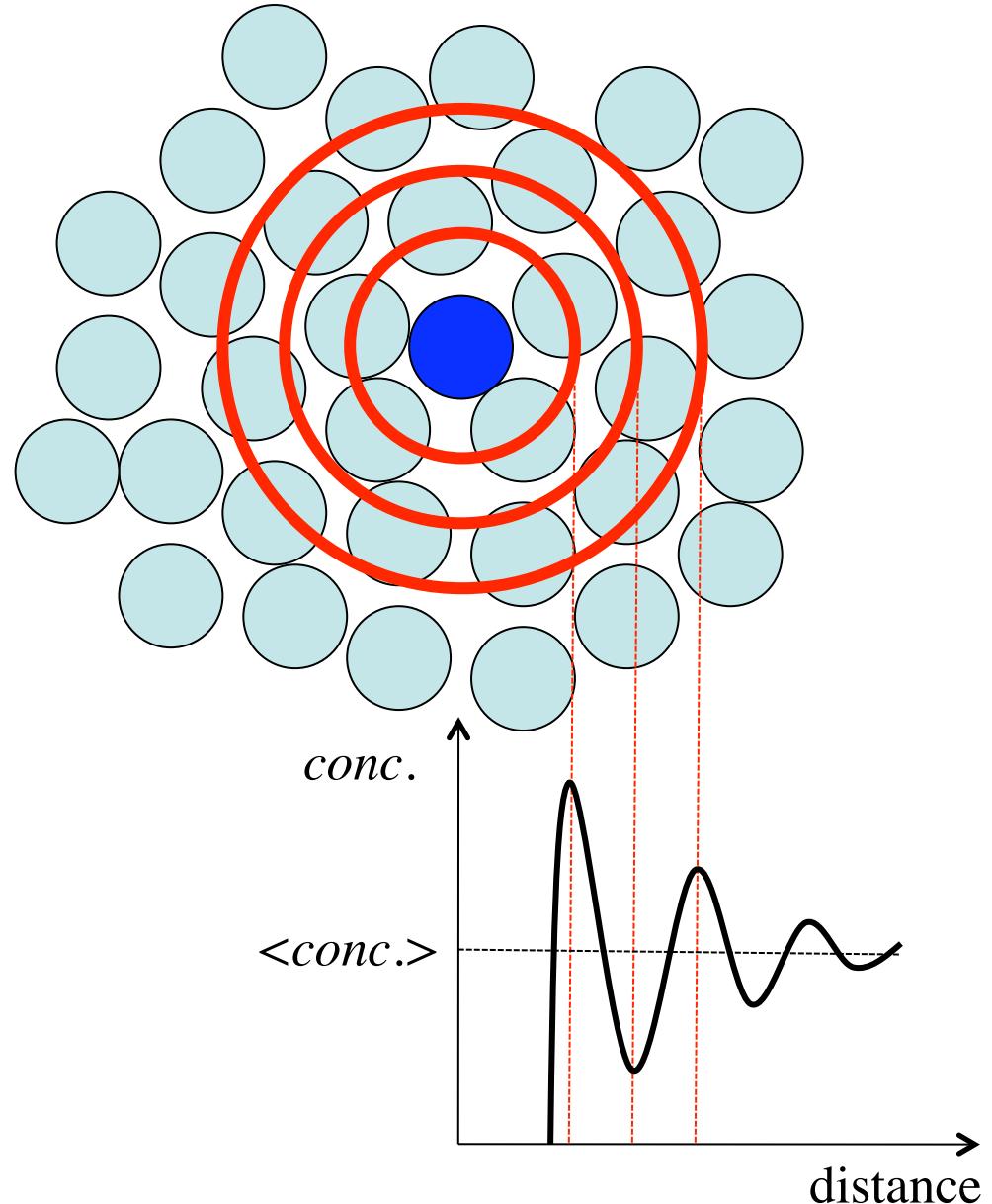
For the analytical expression see e.g.
Kinnig and Thomas *Macromolecules* 17, 1712 (1984)



$S(q) \rightarrow 1$ when $\phi \rightarrow 1$.
i.e. $I(q) \approx P(q)$ at low ϕ .

$S(q) \rightarrow 1$ when $q \rightarrow \infty$.
i.e. $I(q) \approx P(q)$ at high q .

The radial distribution function, $g(r)$



Structure factor

$$S(q) = 1 + \frac{1}{N} \left\langle \sum_m \sum_{n \neq m} e^{i \vec{q} \cdot \vec{r}_{mn}} \right\rangle$$

See *e.g.* J. P. Hansen, I. R. McDonald, “Theory of Simple Liquids”, Ch. 5

In terms of the radial distribution function:

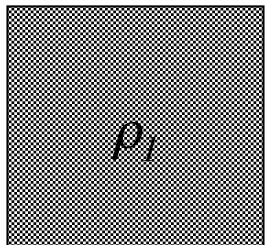
$$S(q) = 1 + \langle c \rangle \int d\vec{r} g(\vec{r}) e^{i \vec{q} \cdot \vec{r}}$$

(or $h(r) = g(r) - 1$)

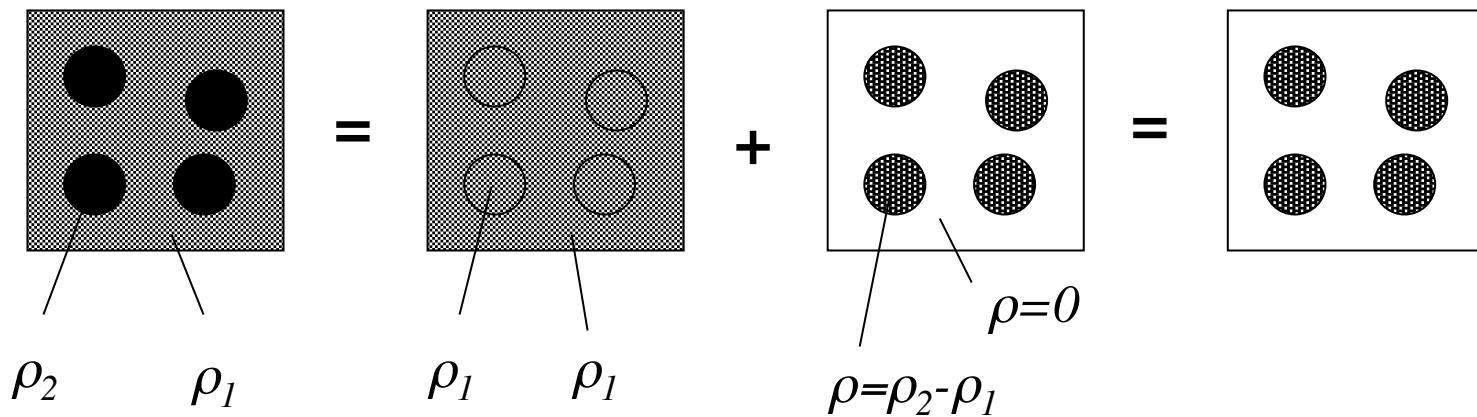
$g(r)$: the *radial distribution function*.

The concentration profile of particles outside (around and away from) a given particle.

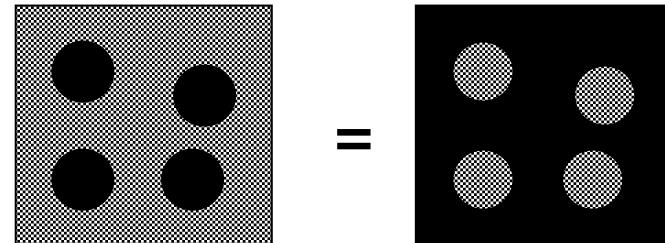
The principle of Babinet



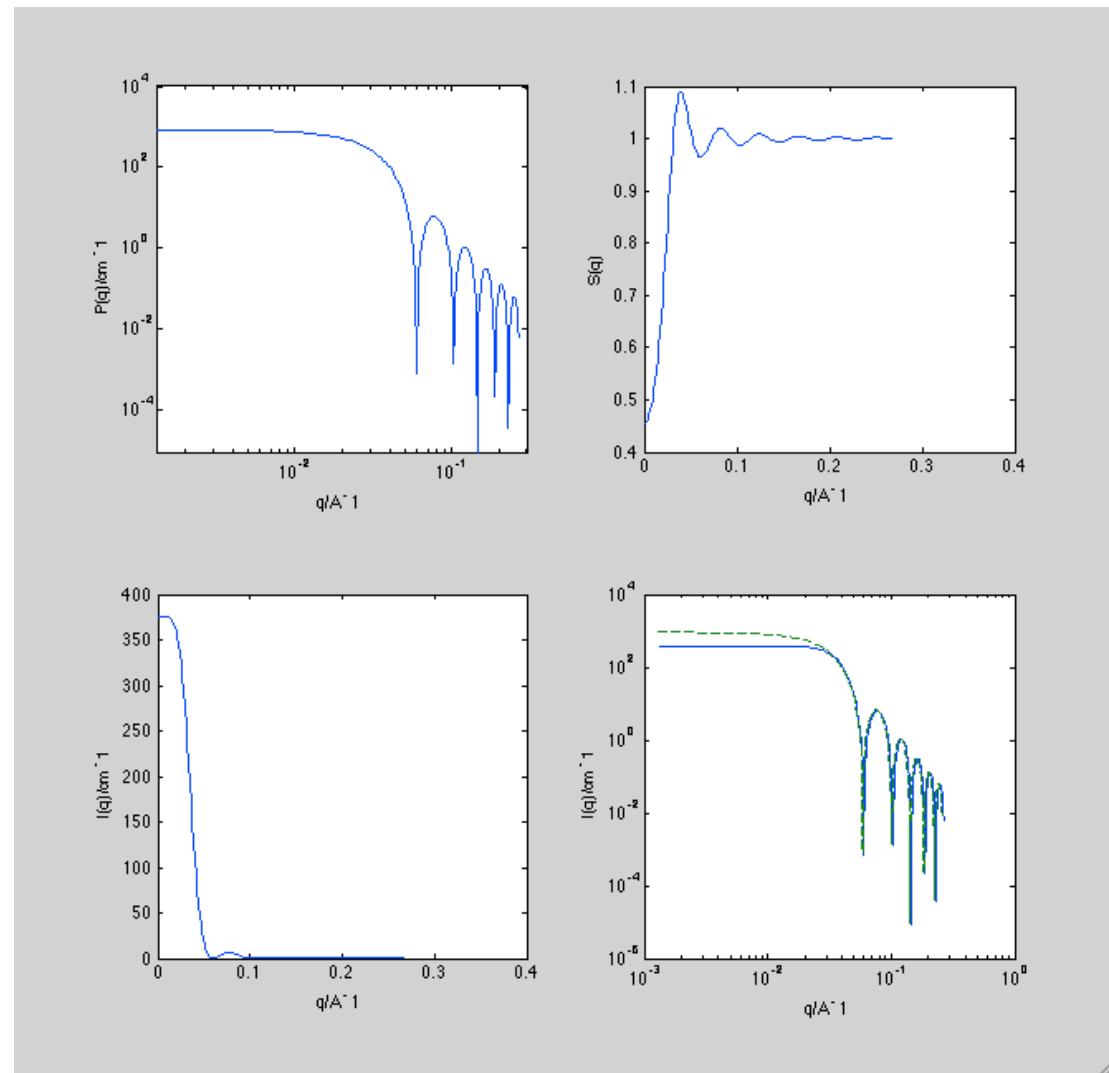
$$I(q) \sim \left(\int_{\mathbb{R}^3} d\vec{r} e^{i\vec{q} \cdot \vec{r}} \right)^2 = \delta(\vec{q})$$



$$I(q) \sim (\Delta\rho)^2$$



$$I(q) = \frac{N}{V} \Delta\rho^2 S(q) P(q)$$



$\phi = 0.1, R = 75 \text{ \AA}, \Delta\rho = 6.83 \text{ } 10^{10} \text{ cm}^{-2}$